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# Measurement of the compressive crack resistance curve of composites using the size effect law 

G. Catalanotti ${ }^{\text {a, }}$, J. Xavier ${ }^{\text {b }}$, P.P. Camanho ${ }^{\text {a }}$<br>${ }^{\text {a }}$ DEMec, Faculdade de Engenharia, Universidade do Porto, Rua Dr. Roberto Frias, 4200-465, Porto, Portugal<br>${ }^{\mathrm{b}}$ CITAB, University of Trás-os-Montes e Alto Douro, Apartado 1013, 5001-801<br>Vila Real, Portugal.


#### Abstract

This paper presents a new methodology to measure the compressive crack resistance curve of the longitudinal plies of carbon-epoxy laminates. The methodology is based on three main steps: the first one corresponds to the determination of the energy release rate of cross-ply laminates with two edge cracks using a parametric finite element model. The energy release rate is used in the definition of a relation between the crack resistance curve and the size effect. Finally, experimental tests are performed in scaled double-edge notched specimens to quantify the size effect law, thus proving the last piece of information required to define the crack resistance curve. The full crack resistance curve is obtained for the IM7-8552 carbon epoxy composite material. The methodology proposed in this paper overcomes the inherent limitations of the existing test methods, and it serves as the basis for the identification of cohesive laws used in some analysis models.


Key words:
A. Polymer-matrix composites (PMCs), B. Fracture toughness, C. Analytical modelling, D. Mechanical testing

## 1 Introduction

The development of a new generation of tougher composite materials presents several challenges to the existing methods for the measurement of the fracture toughness associated with cracks that propagate perpendicularly to the fiber direction. The intra-laminar fracture toughness is relevant not only for material screening and qualification, but also to define the softening laws used in recent analysis models that predict the ultimate strength of composite structures [1-5].

The Compact Tension (CT) test specimen [6,7] is normally used to measure the fracture toughness and the crack resistance curve (R-curve) of composite materials reinforced by unidirectional fibres. While reliable results can be obtained for brittle material systems using appropriate data reduction methods $[8,7]$, the introduction of tougher resins leads to higher loads for crack propagation, which may cause buckling of the unnotched end of the CT test specimen [9].

There is also the need to measure the fracture toughness and the corresponding R-curve associated to the propagation of a kink-band, which shows a crack-like behaviour [10-12] with an R-curve that results from the broadening of the damage height [11]. It is considered here that the compact compression test specimen is inadequate to measure the compressive crack resistance curves of polymer composite materials. In fact, the correction factor used in the data reduction method of the compact compression test method to calculate the energy release rate is the same as that used in the compact tension

[^0]method. However, the contact tractions that occur on the crack faces during a compact compression test render the data reduction method inaccurate. This was demonstrated in a previous investigation [7] where the $\mathcal{J}$-integral around the crack tip was computed using digital image correlation. It was shown [7] that the R -curve of the compact test specimens using the $\mathcal{J}$-Integral and the data reduction method proposed by Pinho et al. [6] are virtually the same for tension but not for compression.

It should also be noted that the compact compression specimen triggers diffused damage during the propagation of the kink-band, artificially increasing the value of the measured fracture toughness, and that it is not possible to identify the location of the tip of the kink band [7]. Therefore, it is considered that while the compact compression test method may be used to measure the initial value of the fracture toughness it does not provide reliable information for the generation of the R -curve.

This means that the analysts have no reliable test methods to measure some of the required material properties, namely the fracture toughness related with the propagation of a kink-band and the corresponding R-curve.

This fact provides the motivation for this paper, whose objective is to propose a new methodology to obtain the R-curve of composite materials that fail by the propagation of a kink-band. The approach proposed is based on the relation between the size effect law and the R-curve, and it will be presented in the following order: section 2 describes the analytical model that relates the size effect law obtained in cross-ply laminates with two edge notched to the R-curve. Section 3 presents the details of the compression tests performed in the double-edge notched specimens manufactured using IM7-8552 carbonepoxy. Section 4 combines the analytical model and the test results to fully propagating in the $x$-direction reads [13]:

$$
\begin{equation*}
\mathcal{G}_{I}=\left(s_{11} s_{22} \frac{1+\rho}{2}\right)^{1 / 2} \lambda^{-1 / 4} \mathcal{K}_{I}^{2} \tag{1}
\end{equation*}
$$

${ }_{59}$ where $s_{l m}$ are the components of the compliance matrix computed in the $x$ -
${ }_{60} y$ coordinate system, $\mathcal{K}_{I}$ is the stress intensity factor, and $\lambda, \rho$ are the two 61 dimensionless elastic parameters defined as:

$$
\begin{equation*}
\lambda=\frac{s_{11}}{s_{22}}, \quad \rho=\frac{2 s_{12}+s_{66}}{2 \sqrt{s_{11} s_{22}}} \tag{2}
\end{equation*}
$$

${ }_{63}$ case, $s_{11}=s_{22}, \lambda=1$ and equation (1) reads:

$$
\begin{equation*}
\mathcal{G}_{I}=\frac{1}{E} \sqrt{\frac{1+\rho}{2}} \mathcal{K}_{I}^{2} \tag{3}
\end{equation*}
$$

where $E$ is the laminate Young's modulus along the $x$ (or $y$ ) direction.

The stress intensity factor of the double edge notched specimen shown in Figure 1 is a function of $\rho$, of the remote stress $\sigma$, and of the shape and the size of the specimen:

$$
\begin{equation*}
\mathcal{K}_{I}=\sigma \sqrt{w} \sqrt{\phi(\alpha, \rho)} \tag{4}
\end{equation*}
$$

where $\alpha=a / w$ is the shape-parameter and $\phi(\alpha, \rho)$ is the correction factor for the geometry and orthotropy of the material. Substituting (4) in (3) the energy release rate reads:

$$
\begin{equation*}
\mathcal{G}_{I}=\frac{1}{E} \sqrt{\frac{1+\rho}{2}} \sigma^{2} w \phi(\alpha, \rho)=\frac{1}{4 w E} \sqrt{\frac{1+\rho}{2}}\left(\frac{P}{t}\right)^{2} \phi(\alpha, \rho) \tag{5}
\end{equation*}
$$

where $t$ is the thickness of the specimen.

Following the approach proposed in [7] the dimensionless function $\phi(\alpha, \rho)$ can be defined for the problem under consideration using the Finite Element Method (FEM). For this purpose a parametric Finite Element model was created using Python [14] together with Abaqus 6.8-3 Finite Element code [15].

In this model the characteristic distance $w$ is taken constant and equal to the unity, while the variables are: i) the shape parameter or, in other words, the crack length $0<\alpha<1$; ii) the dimensionless parameter $\rho$ that takes into account the effect of the orthotropy of the material $(0 \leq \rho \leq 20)$. There is no need to take into account $\lambda$ because for the layup that was chosen (balanced cross ply) its value is constant and equal to one; therefore $\lambda$ is not accounted for in the calibration of the model. Figure 2 shows the mesh of the finite

Using the results obtained in the FE analysis the correction factor $\phi$ is approximated by the following polynomial:

$$
\begin{equation*}
\phi(\alpha, \rho)=\frac{\alpha}{1-\alpha} \sum_{i=1}^{M} \sum_{j=1}^{N} \Phi_{i j} \rho^{j-1} \alpha^{i-1} \tag{6}
\end{equation*}
$$

${ }_{95}$ where $\Phi_{i j}$ is the element of the matrix $\boldsymbol{\Phi}$ at the row $i$ and at the column $j$, ${ }_{96}$ and $M$ and $N$ are the number of rows and columns of $\boldsymbol{\Phi}$ respectively. The 97 matrix $\boldsymbol{\Phi}$ reads:

$$
\boldsymbol{\Phi}=\left[\begin{array}{cccc}
4.315050777 & -0.1833177904 & 0.01642021976 & -4.829962430 \mathrm{E}-4  \tag{7}\\
-5.148136502 & -0.3554678337 & -9.974634025 \mathrm{E}-4 & 4.975387379 \mathrm{E}-4 \\
2.385888075 & 1.339974300 & -0.05966399650 & 7.544565390 \mathrm{E}-4 \\
-0.2810124370 & -0.8040552990 & 0.04491874691 & -7.869467548 \mathrm{E}-4
\end{array}\right]
$$

Figure 3 shows the numerical points obtained in the FE simulations (each point corresponds to one simulation) and the surface fitting function of equation (6).
[Fig. 3 about here.]

Figure 4 shows the relation between $\phi$ and $\alpha$ for different values of the dimensionless elastic parameter $\rho$ obtained by the polynomial approximation and by the FE model. It can be noticed that the fitting obtained is excellent. The curve with $\rho=1$ corresponds to the isotropic case. The close range $\rho \in[0,20]$ should cover all the practical applications. For example for a IM7$8552[0 / 90]_{n s}$ laminate $\left(E_{1}=171420 \mathrm{MPa}, E_{2}=9080 \mathrm{MPa}, G_{12}=5290 \mathrm{MPa}\right.$, $\left.\nu_{12}=0.32\right)$ the elastic parameters calculated using the classical lamination theory are $E_{x}=E_{y}=90648 \mathrm{MPa}, G_{x y}=5290 \mathrm{MPa}, \nu_{x y}=\nu_{x y}=0.032$ and, from (2), $\rho=8.54$.

$$
\text { [Fig. } 4 \text { about here.] }
$$

After defining $\phi$, it is observed that the energy release rate $\mathcal{G}(a)$ is an increasing function of the crack length, or, in other words, the proposed specimen has a positive geometry [17]. Equation (5) can be re-written as:

$$
\begin{equation*}
\mathcal{G}_{I}(a+\Delta a)=\frac{1}{4 w E} \sqrt{\frac{1+\rho}{2}}\left(\frac{P}{t}\right)^{2} \phi\left(\alpha_{0}+\frac{\Delta a}{w}, \rho\right) \tag{8}
\end{equation*}
$$

where $\alpha_{0}=a_{0} / w$ (see Fig. 1). The energy release rate $\mathcal{G}_{I}(a+\Delta a)$ and the schematic representation of the R-curve of the material $\mathcal{R}_{I}(\Delta a)$ are shown in Fig. 5. The dotted line represents the $\mathcal{G}_{I}(a+\Delta a)$ at a constant load $P$. The crack cannot propagate at a constant load if $\mathcal{G}_{I}<\mathcal{R}$ while it will propagate dynamically if $\mathcal{G}_{I}>\mathcal{R}$. The dashed line in Fig. 1, tangent at the R-curve, represents the crack-driving force at the peak load, $P_{u}$ (or at the maximum remote stress, $\left.\sigma_{u}\right)$.

## [Fig. 5 about here.]

In summary, for different sizes $w_{k}$ the driving-force curves $\mathcal{G}_{I}$ corresponding at the peak loads $P_{u k}$ are tangent to R-curve, $\mathcal{R}$. This fact will be used to measure the R -curve.
2.1 Obtaining the $R$-curve from size effect

Based on the previous observations, the peak load, $P_{u}$, or the ultimate nominal stress, $\sigma_{u}=P_{u} /(2 w t)$, can be obtained from the following system of equations:

$$
\left\{\begin{array}{l}
\mathcal{G}_{I}(a+\Delta a)=\mathcal{R}(\Delta a)  \tag{9}\\
\frac{\partial \mathcal{G}_{I}(a+\Delta a)}{\partial \Delta a}=\frac{\partial \mathcal{R}(\Delta a)}{\partial \Delta a}
\end{array}\right.
$$

Assuming that the size effect law, $\sigma_{u}=\sigma_{u}(w)$, is known, substituting (5) in the first of equations (9) yields:

$$
\begin{equation*}
\frac{1}{E} \sqrt{\frac{1+\rho}{2}} \sigma_{u}^{2} w \phi\left(\alpha_{0}+\frac{\Delta a}{w}\right)=\mathcal{R}(\Delta a) \tag{10}
\end{equation*}
$$

131 This equation holds for every $w$. Following [17], differentiating (10) with re132 spect to $w$, under the hypothesis that geometrically similar specimens are ${ }_{133}$ tested ( $\alpha_{0}$ is not a function of $w$ ) and remembering that, by definition, the

$$
\begin{equation*}
\frac{1}{E} \sqrt{\frac{1+\rho}{2}} \frac{\partial}{\partial w}\left(\sigma_{u}^{2} w \phi\left(\alpha_{0}+\frac{\Delta a}{w}\right)\right)=0 \tag{11}
\end{equation*}
$$

$$
\begin{equation*}
\mathcal{R} h \mathrm{~d} a=\mathcal{R}_{0} h_{0} \mathrm{~d} a+\mathcal{R}_{90} h_{90} \mathrm{~d} a \tag{12}
\end{equation*}
$$

where $\mathcal{R}_{0}$ and $\mathcal{R}_{90}$ are the R -curves for the $0^{\circ}$ and $90^{\circ}$ plies respectively while $h, h_{0}$ and $h_{90}$ are respectively the thickness of the laminate, of the $0^{\circ}$ and of the $90^{\circ}$ plies respectively. Taking into account that $\mathcal{R}_{90} \ll \mathcal{R}_{0}$, and $h_{0}=$ $h_{90}=h / 2$, equation (12) results in:

$$
\begin{equation*}
\mathcal{R}_{0}=2 \mathcal{R} \tag{13}
\end{equation*}
$$

Therefore the fracture toughness in longitudinal direction $\mathcal{R}_{0}$ is simply obtained as twice the value of the fracture toughness of the laminate $\mathcal{R}$.

## 3 Experiments

### 3.1 Material and specimen configurations

The material used is the Hexcel IM7-8552 carbon fiber reinforced epoxy. The elastic properties of the material were previously obtained in [18] and are shown in Table 1, where $E_{1}$ and $E_{2}$ are the longitudinal and the transverse Young's modulus respectively, $G_{12}$ is the shear modulus, and $\nu_{12}$ is the major Poisson's ratio.

$$
\text { [Table } 1 \text { about here.] }
$$

The $[0 / 90]_{8 s}$ layup with a nominal thickness of 4 mm is used. The laminate was cured in a hot-press according to the manufacturer's specifications. The curing cycle consists of three different stages: i) keeping the laminate at $110^{\circ} \mathrm{C}$ for 1 hour; ii) increasing the temperature up to $180^{\circ} \mathrm{C}$ and keeping this temperature for 2 hours iii) cooling at $3^{\circ} \mathrm{C} / \mathrm{min}$. The pressure of 7 bar was used during all the curing cycle.

After curing, the laminate was cut using a diamond saw disk to the nominal specimens size and the notches were machined using a vertical mill equipped with 1 mm diameter drill bit.

Six sizes were chosen for the specimens, corresponding to the references A to F. The nominal dimensions of the specimens are shown in Figure 6. Three specimens were tested for each specimen type. It should be noticed that the
initial crack length $a_{0}$ was chosen to be one-half of the characteristic length for all specimens. Therefore the parameter $\alpha_{0}$ is equal to 0.5 .

$$
\text { [Fig. } 6 \text { about here.] }
$$

The specimens manufactured are scaled from the specimen shown in Figure 1. It should be noted that in the actual specimen the crack faces are separated by a finite distance of 1 mm . This ensures that there is no contact of the crack faces; such contact would have invalidated the experimental results. It should also be noted that the shape of the crack tip (semicircular, 1 mm of diameter) does not have an influence on the correct determination of the R -curve of the material because: i) the crack tip shape does not influence the value of the fracture toughness in compression [19], and ii) after the initial propagation the crack tip can be considered sharp and this should not influence the value of the fracture toughness at unstable crack propagation.

After manufacturing, the specimens were painted with a matte white. After drying the specimens, a speckle was made for all the specimens with the help of an airbrush (see Fig. 7).
[Fig. 7 about here.]

### 3.2 Photo-mechanical compression tests

The compression tests were performed with the support of the digital image correlation (DIC) system (Fig. 8). The tests were carried out on a Instron 4208 universal testing machine under displacement control, with a cross-head displacement rate of $0.1 \mathrm{~mm} / \mathrm{min}$ and at room temperature of $22^{\circ} \mathrm{C}$ and relative humidity of $56 \%$. The load was measured using a 100 kN load cell. The

Spider data acquisition system along with Catman software was used to record the strain gauge signal (in the specimens equipped with a strain gauge). 350 $\Omega$ strain gauges from Vishay Micro-Measurement, C2A-06-062LW-350, were used to verify if buckling occurs in the longest specimens (F specimens). To replicate the same frictionless condition imposed numerically for the calculation of the correction factor $\phi$, a thin layer of molybdenum disulfide (MoS2) was used between the specimen and the loading system. The self alignment system shown in Figure 8(b) was also used. The system consists of two parts connected by a spherical joint. The upper part of the self alignment system is in contact with the specimen and it is equipped with a tungsten-carbide insert to prevent indentation; the bottom part is bolted to the frame of the testing machine. By allowing small rotations between the two parts, the self alignment device avoids premature failure of the specimen that may result from possible parallelism errors between the two faces in contact with the loading system.
[Fig. 8 about here.]

At the scale of observation used in this work, the textured pattern required for the digital image correlation was created across the gauge section by means of aerosol and airbrush painting, as shown in Fig. 9. The Aramis DIC-2D v6.0.26 system was used in this work [20,21]. An 8-bit Baumer Optronic FWX20 camera coupled with a Nikon AF Micro-Nikkor 200 mm f/4D IF-ED lens were used for image acquisition. The optical devices and measuring parameters are summarised in Table 2 for each specimen configuration (see also Fig 7).
[Fig. 9 about here.]

A subset size and a subset step of $15 \times 15$ pixel $^{2}$ and $13 \times 13$ pixel $^{2}$ were chosen, respectively (Table 2). In order to guarantee a suitable speckle pattern for this
virtual gauge an airbrush was used (Iwata Custom 181 Micron CM-B model with a fluid nozzle of 0.18 mm in diameter and spray 182 range lower than 60 microns). With this marking technique the accurate uniformity, isotropy and spot size can be achieved. A slightly overlapping of 2 pixels between adjacent subsets was chosen in order to enhance spatial resolution in the strain reconstruction. A strain base length of 5 subsets was selected for strain computation. The displacement and strain resolution of the measured is expected in the range of $2 \times 10^{-2}$ pixel and $0.02-0.04 \%$ respectively (Table 2 ).
[Table 2 about here.]

### 3.3 Experimental results

Representative load vs. displacement curves are shown in Figure 10 for the different specimens.
[Fig. 10 about here.]

The F specimens (the longest specimens) were equipped with a strain gauge on the back side of the specimen (the side without the speckle for the DIC). The strain measured by the strain gauge was compared to the strain measured by the DIC (virtual strain gauge) on the front side in the location corresponding to the physical strain gauge. The strain of the virtual strain gauge was obtained as the average of the longitudinal strain on the rectangular region symmetric to the strain gauge and with the dimension of the electrical grid of the physical strain gauge. Figure 11 shows the comparison between the signal of the physical and of the virtual strain gauge for a representative specimen and the location of the strain gauge in the specimen. As virtually there are no
differences, it is concluded that no buckling occurred in the test specimens.

$$
\text { [Fig. } 11 \text { about here.] }
$$

The digital image correlation can be used to further verify the validity of the tests performed. Figure 12 shows the contour plot of the shear strain, $\gamma_{x y}$. The perfect antisymmetry of the strain field with respect to the longitudinal axis of the specimen and the zero value of the shear strain all over the central part of the specimen (and in particular near the crack tip) shows that the loading system used is appropriate. Any misalignment would result in a different field of the shear strain, and mixed mode crack propagation would have occurred.
[Fig. 12 about here.]

The fracture surface for the two types of specimens (A and F) are shown in Figure 13. It is observed in both specimens that the compressive fracture occurred along the direction of the initial notches.
[Fig. 13 about here.]

The summary of the results is shown in Table 3. For each type of specimen the characteristic size $w$, the peak load $P_{u}$, and the standard deviation are reported. The results obtained are used in the following section in the calculation of the R-curve.

$$
\text { [Table } 3 \text { about here.] }
$$

## 4 Obtaining the R-curve from size effect

As previously explained, the use of the size effect method to obtain the Rcurve of the material requires the knowledge of the size-effect law $\sigma_{u}=\sigma_{u}(w)$ that relates the nominal remote stress to the characteristic size of the body. According to Bažant and Planas [17] different kind of regression plots can be used to find the size-effect laws: i) bilogarithmic plot ( $\ln \sigma_{u}$ vs. $\left.\ln w\right)$; ii) linear regression I ( $\sigma^{-2}$ vs. $w$ ); iii) linear regression II $\left((\sigma w)^{-2}\right.$ vs. $\left.1 / w\right)$; iv) inverse bilogaritmic plot $\left(\ln w\right.$ vs. $\left.\ln \sigma_{u}\right)$. Those methods yield approximatively the same results.

After the inspection of the data obtained experimentally and reported in Table 3, the linear regression I was selected. Figure 14 shows $\sigma^{-2}$ as a function of the characteristic size $w$. The experimental points and the linear fitting are also reported.

$$
\text { [Fig. } 14 \text { about here.] }
$$

The size effect law of the material can be expressed as:

$$
\begin{equation*}
\sigma_{u}=(m w+q)^{-\frac{1}{2}} \tag{14}
\end{equation*}
$$

where $m$ and $q$ are the slope and the intercept of the linear fit respectively. It should be noted that the linear fit was obtained with a coefficient of determination of $R^{2}=0.92$.

Knowing the size effect law, equation (11) can be solved for $w=w(\Delta a)$. All the other parameters are known: $\alpha_{0}=0.5, E=90648 \mathrm{MPa}$, and $\rho=8.54$. Substituting $w=w(\Delta a)$ in equation (10) enables the calculation of the R-
curve of the laminate, $\mathcal{R}$. Finally $\mathcal{R}_{0}$ is obtained by multiplying by two the value of $\mathcal{R}$ for every $(\Delta a)$.

Figure 15 shows the $\mathcal{R}_{0}$ obtained using this methodology as the envelope of the crack driving force curves. The value of the steady-state value of the fracture toughness can be obtained as:

$$
\begin{equation*}
\mathcal{R}_{s s}=\lim _{w \rightarrow \infty} \mathcal{R}_{0}=\frac{\sqrt{2(1+\rho)}}{E} \frac{\phi_{0}}{m}=61 \mathrm{~kJ} / \mathrm{m}^{2} \tag{15}
\end{equation*}
$$

where $\phi_{0}=\left.\phi\right|_{\alpha=\alpha_{0}}$. It is interesting to note that the value of the fracture toughness previously measured using the compact compression test specimens [7], $47.5 \mathrm{~kJ} / \mathrm{m}^{2}$, corresponds to just one point in the rising R-curve.

The length of fracture process zone reads [17]:

$$
\begin{equation*}
l_{f p z}=\frac{f_{0}}{2 f_{0}^{\prime}} w_{0} \tag{16}
\end{equation*}
$$

where $f_{0}=\left.\sqrt{\phi}\right|_{\alpha=\alpha_{0}}, f_{0}^{\prime}=\partial \sqrt{\phi} /\left.\partial \alpha\right|_{\alpha=\alpha_{0}}$ and $w_{0}$ is the constant of Bažant sizeeffect law with the dimension of a length. For the linear regression used [17] $w_{0}=q / m ;$ therefore equation (16) reads:

$$
\begin{equation*}
l_{f p z}=\frac{f_{0}}{2 f_{0}^{\prime}} \frac{q}{m}=1.43 \mathrm{~mm} \tag{17}
\end{equation*}
$$

[Fig. 15 about here.]

To simplify the use of the R-curve it is necessary to express it in an analytical form that fits the points obtained by solving equations (10) and (11). The following expression can be used [17]:

$$
\begin{cases}\mathcal{R}_{0}=\mathcal{R}_{s s}\left[1-(1-\kappa \Delta a)^{n}\right] & \text { if } \Delta a<l_{F P Z}  \tag{18}\\ \mathcal{R}_{0}=\mathcal{R}_{s s} & \text { if } \Delta a \geq l_{F P Z}\end{cases}
$$

By optimal fitting the following values are obtained: $\kappa=0.5126$ and $n=4.289$. Figure 16 shows the analytical fitting of the R -curve.
[Fig. 16 about here.]

## 5 Conclusion

This paper demonstrates that it is possible to obtain the compressive R-curve of the longitudinal plies of carbon-epoxy laminates by quantifying the size effect law observed in compressive tests of scaled specimens with two edge cracks. The limiting factors of the previously proposed compact compression test method, such as the difficulty in predicting the exact location of the extremity of the kink band, do not play any role in the model developed in this paper.

Using both electrical strain gauges and the digital image correlation system it is demonstrated that the test method and specimens' geometry proposed do not result in premature buckling failure. In addition, the strain field measured by digital image correlation system indicates that the specimens are properly aligned. Therefore, the peak loads measured in the tests can be confidently used for the definition of the size effect law.

While the previous tests methods are only able to provide one single value for the compressive fracture toughness, the proposed method enables the identification of the full R-curve. The steady-state value of the fracture toughness
of the IM7-8552 carbon-epoxy composite material is $61 \mathrm{~kJ} / \mathrm{m}^{2}$, corresponding to a length of the fracture process zone of 1.43 mm .

The work presented in this paper can be used as the basis for material screening and selection, to identify cohesive laws used in the recent analysis models, and to estimate the fracture toughness of multidirectional laminates using the fracture toughness of the $0^{\circ}$ plies.

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Fig. 1. Double edge notched compression specimen.

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Fig. 2. Finite element model of the specimen.


Fig. 3. Relation between $\phi$ and $\alpha$ for different values of $\rho$.


Fig. 4. Numerical points and $\phi$ for different values of $\rho$.


Fig. 5. Crack driving force and resistance curves.


Fig. 6. Specimens' geometry (dimensions in mm ).


Fig. 7. Specimens' configuration.

(a) Optical set-up

(b) Loading system

Fig. 8. Photomechanical set-up.


Fig. 9. Speckle pattern typically used in the digital image correlation measurements (region of interest of $1543 \times 1025$ pixel $^{2}\left(22.0 \times 14.6 \mathrm{~mm}^{2}\right)$ ).


Fig. 10. Typical load displacement curves for the specimens A, B, C, D and F.


Fig. 11. Comparison between the signal of the physical (SG) and virtual (DIC) strain gauge.

1


Fig. 12. Typical contour plot of the shear strain, $\gamma_{x y}$.

(a) type A

(b) type F

Fig. 13. Tested specimens (not at scale).


Fig. 14. $\sigma_{u}^{-2}$ vs. $w$ and linear fitting.


Fig. 15. Envelope of driving force curves and R-curve.


Fig. 16. Analytical fitting of the R-curve.

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| $E_{1}[\mathrm{MPa}]$ | $E_{2}[\mathrm{MPa}]$ | $G_{12}[\mathrm{MPa}]$ | $\nu_{12}$ |
| :---: | :---: | :---: | :---: |
| 171420 | 9080 | 5290 | 0.32 |

Table 1
Elastic properties of IM7-8552.

Table 2
Optical system components and measurement parameters (see Fig. 7 for A, B, C, $\underline{\mathrm{D}, \mathrm{E} \text { and } \mathrm{F} \text { specimen configuration). }}$

| CCD camera | Baumer Optronic FWX20 |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | ( 8 bit, $1624 \times 1236$ pixels) |  |  |  |  |  |
| Lens | (Working F-number: 11) |  |  |  |  |  |
| Lighting system | LED RAYL | X 25 |  |  |  |  |
| Image recording | A | B | C | D | E | F |
| Field Of View $\left(\mathrm{mm}^{2}\right)$ | $15.2 \times 11.5$ | $23.2 \times 17.7$ | $31.2 \times 23.7$ | $38.8 \times 29.5$ | $46.2 \times 35.1$ | $55.3 \times 42.1$ |
| Working Distance |  |  |  |  |  |  |
| Conversion factor (mm/pixel) | 0.0093 | 0.0143 | 0.0192 | 0.0239 | 0.0284 | 0.034 |
| Acquisition frequency |  |  |  |  |  |  |
| (Hz) | 1 | 1 | 1 | 0.5 | 0.5 | 0.5 |
| Exposure time (ms) | 4 | 4 | 7 | 7 | 12 | 20 |
| DIC (Aramis) parameters |  |  |  |  |  |  |
| Subset size | $15 \times 15$ pixe |  |  |  |  |  |
| Subset step | $13 \times 13$ pixe |  |  |  |  |  |
| Strain base length | $5 \times 5$ subset |  |  |  |  |  |
| DIC resolution |  |  |  |  |  |  |
| Displacement | $2 \times 10^{-2}$ pix |  |  |  |  |  |
| Strain | 0.02-0.04\% |  |  |  |  |  |


| specimen label | A | B | C | D | E | F |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: |
| $w[\mathrm{~mm}]$ | 5 | 7.5 | 10 | 12.5 | 15 | 17.5 |
| $P_{u}[\mathrm{~N}]$ | 11390 | 15477 | 17436 | 22622 | 24706 | 25628 |
| $\operatorname{STDV}\left(P_{u}\right)[\mathrm{N}]$ | 2262 | 1421 | 3268 | 1793 | 1827 | 1028 |
| $\sigma_{u}[\mathrm{MPa}]$ | 285 | 258 | 218 | 226 | 206 | 183 |

Table $\overline{\overline{3}}$
Summary of the experimental results.


[^0]:    * Corresponding author

    Email address: giuseppe.catalanotti@fe.up.pt (G. Catalanotti).

